

# Seismic reflection images of a near-axis melt sill within the lower crust at the Juan de Fuca ridge

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The oceanic crust extends over two-thirds of the Earth's solid surface, and is generated along mid-ocean ridges from melts derived from the upwelling mantle<sup>1</sup>. The upper and middle crust are constructed by dyking and sea-floor eruptions originating from magma accumulated in mid-crustal lenses at the spreading axis<sup>2–6</sup>, but the style of accretion of the lower oceanic crust is actively debated<sup>7</sup>. Models based on geological and petrological data from ophiolites propose that the lower oceanic crust is accreted from melt sills intruded at multiple levels between the Moho transition zone (MTZ) and the mid-crustal lens<sup>8–11</sup>, consistent with geophysical studies that suggest the presence of melt within the lower crust<sup>12–16</sup>. However, seismic images of molten sills within the lower crust have been elusive. Until now, only seismic reflections from mid-crustal melt lenses<sup>2,17,18</sup> and sills within the MTZ have been described<sup>19</sup>, suggesting that melt is efficiently transported through the lower crust. Here we report deep crustal seismic reflections off the southern Juan de Fuca ridge that we interpret as originating from a molten sill at present accreting the lower oceanic crust. The sill sits 5–6 km beneath the sea floor and 850–900 m above the MTZ, and is located 1.4–3.2 km off the spreading axis. Our results provide evidence for the existence of low-permeability barriers to melt migration within the lower section of modern oceanic crust forming at intermediate-to-fast spreading rates, as inferred from ophiolite studies<sup>9,10</sup>.

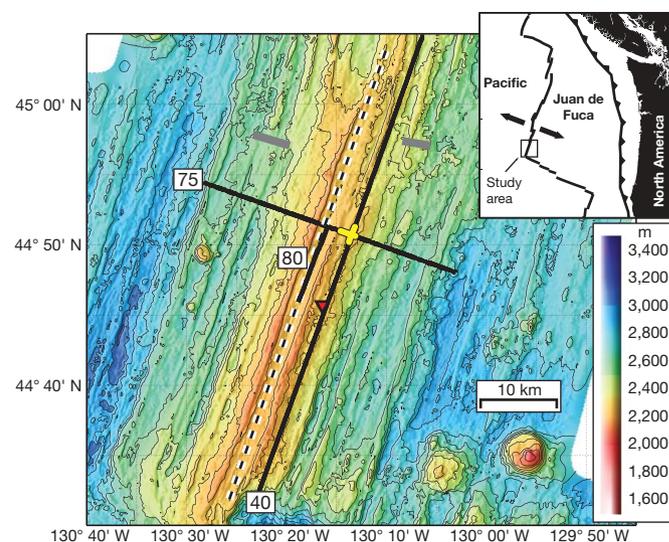
The Juan de Fuca ridge (JdFR) is a mid-ocean ridge in the northeast Pacific (Fig. 1), where spreading occurs at intermediate rates (56 mm yr<sup>-1</sup>; ref. 20). The results presented here arise from a two-dimensional seismic reflection survey conducted in 2002 along the JdFR and its flanks<sup>19,21,22</sup>. This survey imaged the top of an axial magma chamber (AMC) beneath all JdFR segments at depths of ~2–3 km (ref. 22). Reflections from the crust–mantle boundary (the Moho) were observed throughout much of the survey at 2,080 ± 100 ms crustal two-way travel time (TWTT), as well as sub-Moho reflections interpreted as gabbro lenses and possibly melt accumulations<sup>19</sup> (Fig. 1).

Along the Cleft segment of the JdFR (Fig. 1), the top of the AMC is 2.0–2.3 km deep<sup>21</sup> (Fig. 2b, f), and the Moho is imaged intermittently throughout the area<sup>21</sup>. Crustal travel times to the Moho decrease away from the axis and decrease northward, from ~2,550 ms at the south end to ~2,300 ms at the north end<sup>21</sup> (Fig. 2a and Supplementary Fig. 1a). The exception to these general trends is a pronounced ~11-km-long area of anomalously large crustal travel times located ~3–5 km east of the spreading axis (Supplementary Fig. 1a), suggesting local crustal thickening, elevated crustal temperatures, presence of melt above the Moho, or a combination of any of these factors.

At the southern end of this anomaly, profiles 75 and 40 (Fig. 1) show a reflection event at 2,000 ms below the sea floor (Fig. 2a, b, d and h).

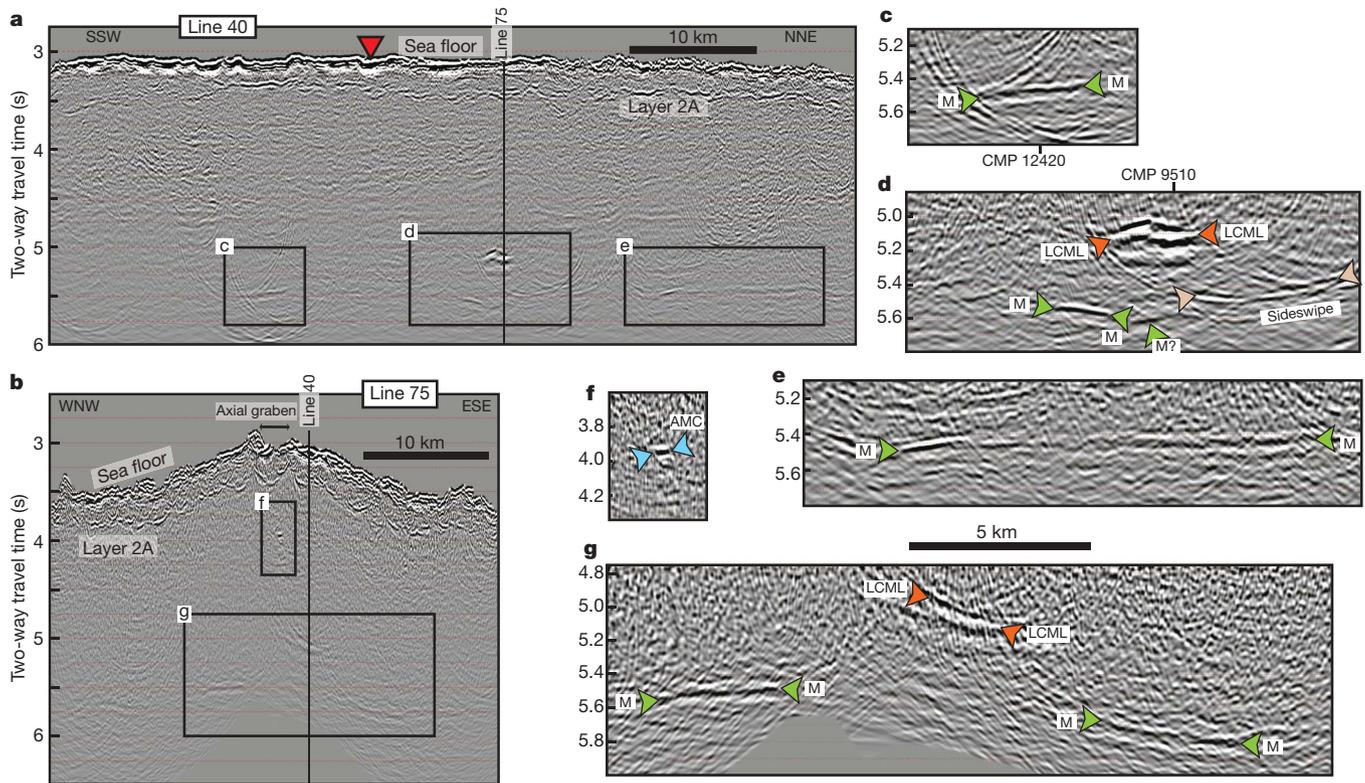
This event is located 1.4–3.2 km to the east of the spreading axis, in 50–114-kyr-old crust<sup>20</sup>, extending ~2.4 km and ~1.8 km in the ridge-parallel and ridge-perpendicular directions, respectively. The reflection event must originate at depths of 5–6 km (for an average crustal velocity of 5–6 km s<sup>-1</sup> above this reflector) from layering that is likely to be at least ~100-m thick in order to be detected by our signal (which has a dominant frequency of 12 Hz and a dominant wavelength of 400–500 m at those depths). The origin of this event could be an anomalous Moho, a frozen sill of ultramafic material embedded in gabbros (such as a wehrlite intrusion<sup>23</sup>), or a molten sill. On the basis of its amplitude variation with offset (AVO) response (Fig. 3) and on other considerations discussed below, we interpret this event as originating from a lower-crustal melt lens (LCML).

The possibility that the reflection event is an anomalously shallow Moho seems unlikely, because this implies abrupt crustal thinning by either a reduction in magma supply in a small area that is otherwise surrounded by large Moho travel times (Supplementary Fig. 1a), or



**Figure 1** | Seismic reflection profiles located over a bathymetric map of the Cleft segment of the southern JdFR. Colours represent sea-floor depth in metres, with contours every 100 m. Numbered solid lines are seismic profiles shown in Figs 2 and 4. Dashed line locates the spreading axis. Yellow lines locate the plan-view extent of the LCML reflection event shown in Figs 2 and 4. Thick grey lines locate sub-Moho reflection events identified<sup>19</sup> in the vicinity of Cleft segment. Red triangle locates low-temperature hydrothermal vents observed on the sea floor<sup>30</sup>. Top-right inset shows location of the study area in the northeast Pacific. Names indicate tectonic plates; arrows indicate direction of sea-floor spreading.

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**Figure 2 | Migrated seismic reflection profiles 40 and 75.** **a**, Ridge-parallel profile 40 is located 3.3 km to the east of the spreading axis (Fig. 1). Vertical line, intersection with profile 75. Boxed areas are shown in detail in **c**, **d** and **e**. Red triangle, low-temperature hydrothermal vents observed on the sea floor<sup>30</sup>. **b**, Ridge-perpendicular profile 75. Vertical line, intersection with profile 40. Spreading axis is within the labelled axial graben. Boxed areas are shown in detail in **f** and **g**. In **c–g**, main events are indicated with coloured arrows: top of axial magma chamber (AMC), lower-crustal melt lens

(LCML) and Moho (M). Vertical axes show two-way travel time in seconds for all panels. Scale bar at lower right is for **c–g**. Locations of CMP 12420 (Fig. 3a) and 9510 (Fig. 3b) are labelled in **c** and **d**, respectively. The event labelled as sideswipe in **d** is interpreted as such based on its steep hyperbolic diffraction character observed in the stacked image (Supplementary Fig. 2a, d), as well as in its higher frequency content relative to the lower-frequency, flat-lying Moho events. Unmigrated stack versions of the images shown here are presented in Supplementary Fig. 2.

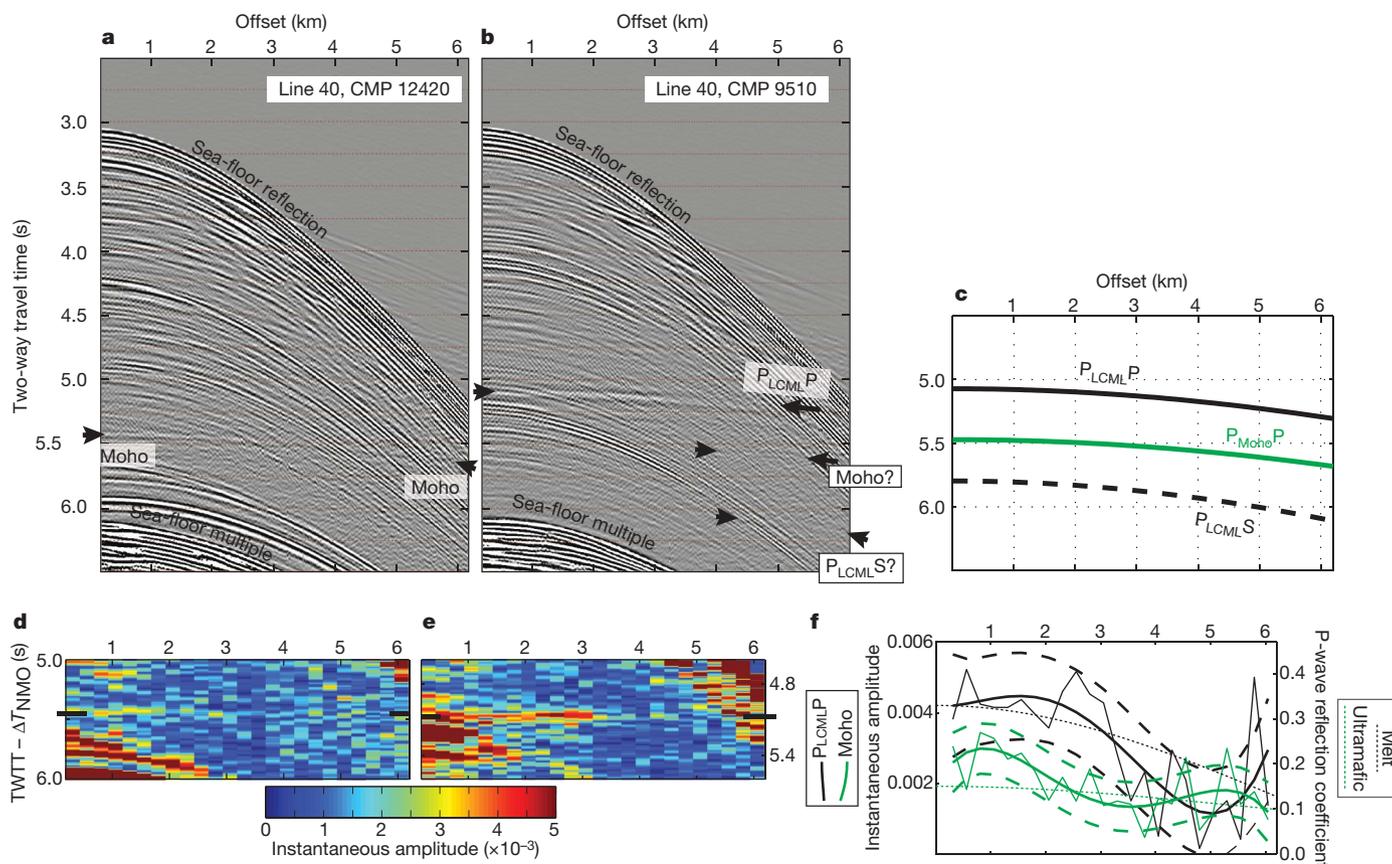
by tectonic thinning or uplifting, for which there is no evidence in the sea-floor morphology. Average crustal TWTT to Moho observed along line 40 is 2,256 ms (Fig. 2a, c–e), which is 256 ms more than for the LCML event (Fig. 2d). In addition, a weak Moho event, which can be recognized in the pre-stack gathers at far offsets (Fig. 3b), appears to extend beneath the LCML (Fig. 2d). This indicates that the top of the LCML is located within the lower crust, about 850–900 m above the Moho (for a lower-crustal velocity of  $6.8 \text{ km s}^{-1}$ ).

The LCML event displays a complex reflection pattern with two positive (black) peaks compared to one from the Moho reflection (Fig. 2d, g). This suggests that the LCML event was produced by a complex velocity structure (rather than a simple velocity increase/decrease), possibly associated with two or more sills or the roof and floor of a single sill. In this case, the polarity of the reflection of the LCML may not be useful to discriminate between positive or negative impedance contrasts, as one would expect from the ideal case of a simple sill intrusion<sup>18</sup>. In a first-order approximation, both the Moho and a frozen ultramafic sill should have identical impedance contrasts because they both represent ultramafic material sitting below gabbros. Therefore, minor amplitude interference effects aside, the comparison of the AVO responses of Moho and LCML events (Fig. 3d–f) observed in pre-stack gathers (Fig. 3a, b) can help determine if the LCML corresponds to a positive or a negative velocity contrast. Figure 3d–f shows that the P-wave reflecting off the LCML ( $P_{\text{LCML}}P$ ) has higher amplitude than the Moho for offsets <3 km, and that the amplitude of the  $P_{\text{LCML}}P$  decays more rapidly between 3 and 5 km offset while the Moho amplitude shows less variation across all offsets (Fig. 3f). The observed AVO response of the LCML resembles the predicted variation with offset of the reflection

coefficient for a P-wave reflecting off a low-velocity molten sill (Fig. 3f). In contrast, the observed AVO response of the Moho event is more consistent with the predicted variation with offset of the reflection coefficient for a P-wave reflecting off a high-velocity Moho or ultramafic sill (Fig. 3f). This suggests that the LCML reflection corresponds to a molten sill.

For a fully molten sill, one expects to observe at far offsets significant converted S-wave energy reflecting off the LCML<sup>18</sup> ( $P_{\text{LCML}}S$ , Supplementary Fig. 3). We do not observe high-amplitude S-wave arrivals in the data, but we recognize a weak arrival at offsets larger than ~4.5 km between ~6.0 and 6.2 s, where the  $P_{\text{LCML}}S$  is expected to be present (Fig. 3b, c). Although the weak amplitude of the  $P_{\text{LCML}}S$  is more consistent with a ‘mush’ sill than a fully molten one (Supplementary Fig. 3), the observed AVO response of the  $P_{\text{LCML}}P$  event suggests a melt-rich sill (other factors such as attenuation may be affecting the strength of the  $P_{\text{LCML}}P$  event). However our data cannot distinguish between mush at the crystal scale and a set of molten sills that are smaller than the dominant wavelength of the seismic signal and encased in solid cumulates.

The largest crustal TWTT to Moho at the southern edge of the LCML event is 2,508 ms, which is 252 ms more than the average crustal TWTT observed along line 40. Assuming that this Moho travel-time anomaly is entirely caused by thickening of the lower crust beneath the LCML, it implies a local thickness anomaly of ~850–900 m (for a lower-crustal velocity of  $6.8 \text{ km s}^{-1}$ ), placing the crust–mantle boundary 1,700–1,800 m beneath the LCML. Alternatively, the Moho travel-time anomaly could be due to a 45–50% reduction in P-wave velocity, implying a completely molten lowermost crust regardless of how melt is distributed and interconnected<sup>24</sup>. These



**Figure 3 | Observed and modelled reflection amplitudes and travel times.** **a, b**, CMP supergathers 12420 (**a**) and 9510 (**b**) from profile 40 (locations in Fig. 2c, d). P waves reflecting off the Moho and the LCML ( $P_{LCML}P$ ) are indicated. Weak but coherent arrivals are observed in CMP 9510 at  $\sim 4.0$ – $5.5$  km offset between  $\sim 5.5$  and  $5.6$  s, and at far offsets ( $>5.5$  km) between  $\sim 6.0$  and  $6.2$  s, and are tentatively interpreted as Moho reflections and S-converted waves reflecting off the LCML ( $P_{LCML}S$ ), respectively. (The  $P_{LCML}S$  would be converted back to P at the sea floor to be detected by the hydrophone streamer.) **c**, Predicted travel-time curves for Moho reflection,  $P_{LCML}P$  and  $P_{LCML}S$  (Methods). **d, e**, Instantaneous amplitude of CMP supergathers 12420 (**d**) and 9510 (**e**) within a 1-s window around the Moho and  $P_{LCML}P$  arrivals, respectively (Methods). Vertical axis shows two-way travel time corrected for normal move-out ( $\Delta T_{NMO}$ ). **f**, AVO analysis of Moho and  $P_{LCML}P$  events (green and black lines, respectively); amplitude

scale at left vertical axis. Thin solid lines, instantaneous amplitudes versus offset at a constant time of 5.445 s for CMP 12420, and 5.075 s for CMP 9510 (short thick lines in **d** and **e**, respectively). Thick solid lines, least-squares cubic fits (dashed lines are 50% confidence limits of the polynomial fits). The increase in instantaneous amplitude for CMP 9510 at offsets  $>5$  km is due to interference of high-amplitude, steeply dipping sea-floor reflections and diffractions with the  $P_{LCML}P$  event, and it does not accurately represent the AVO response of this event. Predicted P-wave reflection coefficients versus source–receiver offset (amplitude scale at right vertical axis) are shown for two scenarios (Methods): the lower medium is a molten sill (short-dashed black line, ‘melt’), or it represents mantle or a solid ultramafic sill (short-dashed green line, ‘ultramafic’). Horizontal axes in all panels are source–receiver offset.

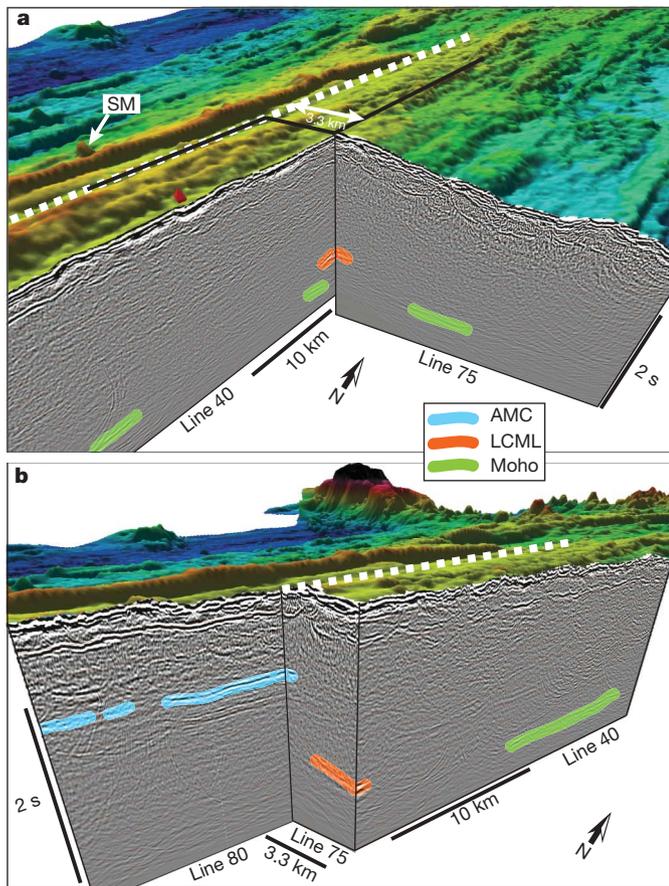
two scenarios are end-member possibilities; most probably the travel-time anomaly below the LCML is due to both presence of melt and local thickening of the lower crust.

Melt migration from the mantle to the crust is thought to be a process of focused porous flow<sup>25,26</sup>, although this may change at the MTZ as melts enter a conductively cooling regime and build up permeability barriers by crystallizing plagioclase, leading to the emplacement of melt sills within the MTZ<sup>11</sup>. In contrast, porous flow may not be the dominant mechanism of melt migration above the MTZ<sup>27</sup>, and lower-crustal sills could be fed by focused flow along channels originating at the MTZ melt sills<sup>10,11</sup>. Several mechanisms have been proposed for the creation of permeability barriers that can allow the emplacement of melt sills at multiple levels within the lower crust, such as the presence of anorthositic bands, plagioclase-clinopyroxene-orthopyroxene saturation within crystallizing gabbros, or formation of low-porosity cumulate gabbros after melt extraction<sup>10</sup>. Our data cannot discriminate between these processes; however, the observation that the LCML is emplaced a few kilometres off the spreading axis allows us to hypothesize that here, the melt sill could have formed at a permeability barrier associated with a contrasting thermal regime above it, in a similar manner to the emplacement of the axial mid-crustal melt lens<sup>28,29</sup>. The axial region of

fast-spreading ridges is characterized by a 5–7-km-wide zone of low seismic velocities<sup>14</sup>. The lateral boundaries of this region are interpreted as steep isotherms controlled by the transition from shallow hydrothermal circulation at the ridge axis to deep-penetrating hydrothermal fluid pathways just a few kilometres off the spreading axis<sup>14</sup>. Thus we interpret the off-axis deep melt sill imaged by our data to represent a deep freezing horizon that roughly follows the boundaries of an axial low-velocity zone.

The imaged LCML is wider than the shallow AMC lens (Fig. 2f, h and Fig. 4). However its along-axis extent is much more limited than the AMC (Fig. 4b), which extends along tens of kilometres beneath the ridge<sup>21</sup>. The along-axis continuity of AMC lenses observed at many spreading centres implies that mid-crustal melt sills are relatively steady-state features frequently replenished with magma from below. Therefore the limited spatial extent of the LCML, together with the absence of similarly sized lower-crustal reflections in the only three-dimensional seismic reflection study of a fast-spreading ridge published to date<sup>12</sup>, suggests that lower-crustal melt sills are probably less frequently fed than AMC lenses, and/or that the LCML imaged in our data is uncommonly large.

The sea floor in the area where profiles 40 and 75 intersect is flatter, smoother, and has higher reflectivity than the surroundings



**Figure 4 | Three-dimensional perspective of crustal structure and sea-floor topography.** **a**, Southern and eastern parts of profiles 40 and 75, respectively, are shown hanging from the sea-floor topography. Moho and LCML events identified in Fig. 2 are emphasized with colour markers. White dashed line, spreading axis. SM, young-looking, near-axis seamount with chemistry distinct from the axial lava flows<sup>30</sup>. Red triangle, low-temperature hydrothermal vents observed on the sea floor<sup>30</sup>. Solid lines, profiles shown in **b**. **b**, Profile 80, northern part of profile 40, and central section of profile 75, all shown hanging from the sea-floor topography. Moho, LCML and AMC events identified in Fig. 2 are emphasized with colour markers. White dashed line, spreading axis, coincident with profile 80.

(Supplementary Fig. 1b–d), suggesting younger terrain and repaving of the sea floor. Lava flows up to 4 km off-axis emanating from ridge-parallel faults and fissures have been identified along the Cleft segment and interpreted as erupting from the edges of the AMC lens<sup>30</sup>. In addition, the axial summit graben of Cleft segment is flanked by some small, young-looking near-axis seamounts with chemistry distinct from the axial lavas (Fig. 4a)<sup>30</sup>. Observations by submersible vehicles have also reported the presence of low-temperature hydrothermal venting ~3.3 km east of the spreading axis<sup>30</sup> and ~10 km to the south of the LCML (Figs 1, 2a, 4a). Our discovery of lower-crustal melt sills at similar distances from the spreading axis in the same area suggests that deep lower-crustal melt lenses could also contribute to construction of the upper crust by off-axis sea-floor eruptions, and provide heat for off-axis hydrothermal circulation.

Our data provide evidence for a molten sill at present accreting the lower oceanic crust at the JdFR. The location of the sill within 50–114-kyr-old crust suggests that a crustal column formed at this section of the JdFR represents magmatic accretion over a time period of at least 100 kyr. Near-axis lower-crustal reflections have been previously reported at the East Pacific Rise<sup>31</sup>. However, their along-axis extent has not been explored and they were not interpreted as melt sills<sup>31</sup>. Our results imply the existence of low-permeability barriers to melt migration within the lower oceanic crust, in agreement with models based on observations at ophiolites<sup>9,10</sup>, and indicate that

melt migration from the MTZ to the shallow AMC lens may not necessarily be as efficient as recently inferred from three-dimensional seismic reflection data<sup>12</sup>. As the palaeo-conditions under which crustal sections exposed at ophiolites were formed are in most cases difficult to determine, there is some ambiguity about the extent to which ophiolites are analogues to modern oceanic crust<sup>32</sup>. Our results thus support the idea that some of the lower-crustal processes inferred from ophiolite settings do occur during the accretion of modern oceanic crust at intermediate-to-fast spreading rates. However, the uniqueness and limited extent of the imaged LCML compared to the mid-crustal melt lens suggests that large lower-crustal melt sills are rare and probably contribute a small fraction to the overall thickness of the lower crust, as suggested by numerical models and petrological constraints<sup>7,33</sup>.

## METHODS SUMMARY

**Data processing.** Acquisition parameters are described in refs 19, 21. Pre-stack processing consisted of *f-k* filtering, sorting into 80-fold common-midpoint (CMP) gathers, amplitude correction for spherical divergence, trace editing, band-pass filtering (2–7–20–40 Hz), velocity analysis and normal move-out correction. After stacking, data were muted below first sea-floor multiple, time migrated using stacking velocities, and muted above the sea-floor reflection. CMP supergathers were constructed by sorting 30 consecutive CMP gathers into fivefold constant-offset gathers, partial stacking, and band-pass filtering (2–7–20–40 Hz). **AVO analysis.** Instantaneous amplitudes were calculated on CMP supergathers after normal move-out correction to flatten the seismic events and partial stacking of traces within 250-m bins.

**Full Methods** and any associated references are available in the online version of the paper at [www.nature.com/nature](http://www.nature.com/nature).

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**Supplementary Information** is linked to the online version of the paper at [www.nature.com/nature](http://www.nature.com/nature).

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## METHODS

**Travel-time curves.** Travel-time curves were calculated assuming a one-dimensional model consisting of a 2,319-m-thick water layer ( $v_p = 1,500 \text{ m s}^{-1}$ ), a 5,940-m-thick crustal layer above the LCML ( $v_p = 6,000 \text{ m s}^{-1}$ ), and a 1,000-m-thick layer between the LCML and the Moho ( $v_p = 5,000 \text{ m s}^{-1}$ ). Crustal Poisson's ratio is 0.25.

**Reflection coefficients.** Reflection coefficients were calculated by assuming two homogeneous, semi-infinite media separated by a flat interface. We considered two models with the same upper medium properties but different lower medium properties. For the upper medium representing the gabbroic lower crust, we use the following elastic parameters:  $v_p = 6,500 \text{ m s}^{-1}$ ,  $v_s = 3,550 \text{ m s}^{-1}$  and  $\rho = 2,800 \text{ kg m}^{-3}$ . These parameters represent average laboratory measurements made at room temperature and 200 MPa confining pressure in gabbros drilled at the Mid-Atlantic Ridge<sup>34</sup> after correcting for the expected elevated temperatures of the near-axis lower crust (900 °C; ref. 7), assuming that  $v_p$  decreases with temperature as  $-0.57 \times 10^{-3} \text{ km s}^{-1} \text{ K}^{-1}$  (ref. 35) and a Poisson's ratio of 0.29 (refs 34, 36).

In the first model, the lower medium represents a melt sill. We use the definition of ref. 18: melt refers to liquid matrix with disconnected crystals, with  $v_s \approx 0 \text{ km s}^{-1}$ . We use the following elastic parameters:  $v_p = 3,400 \text{ m s}^{-1}$  (ref. 18),  $v_s = 0 \text{ m s}^{-1}$  and  $\rho = 2,700 \text{ kg m}^{-3}$  (ref. 28).

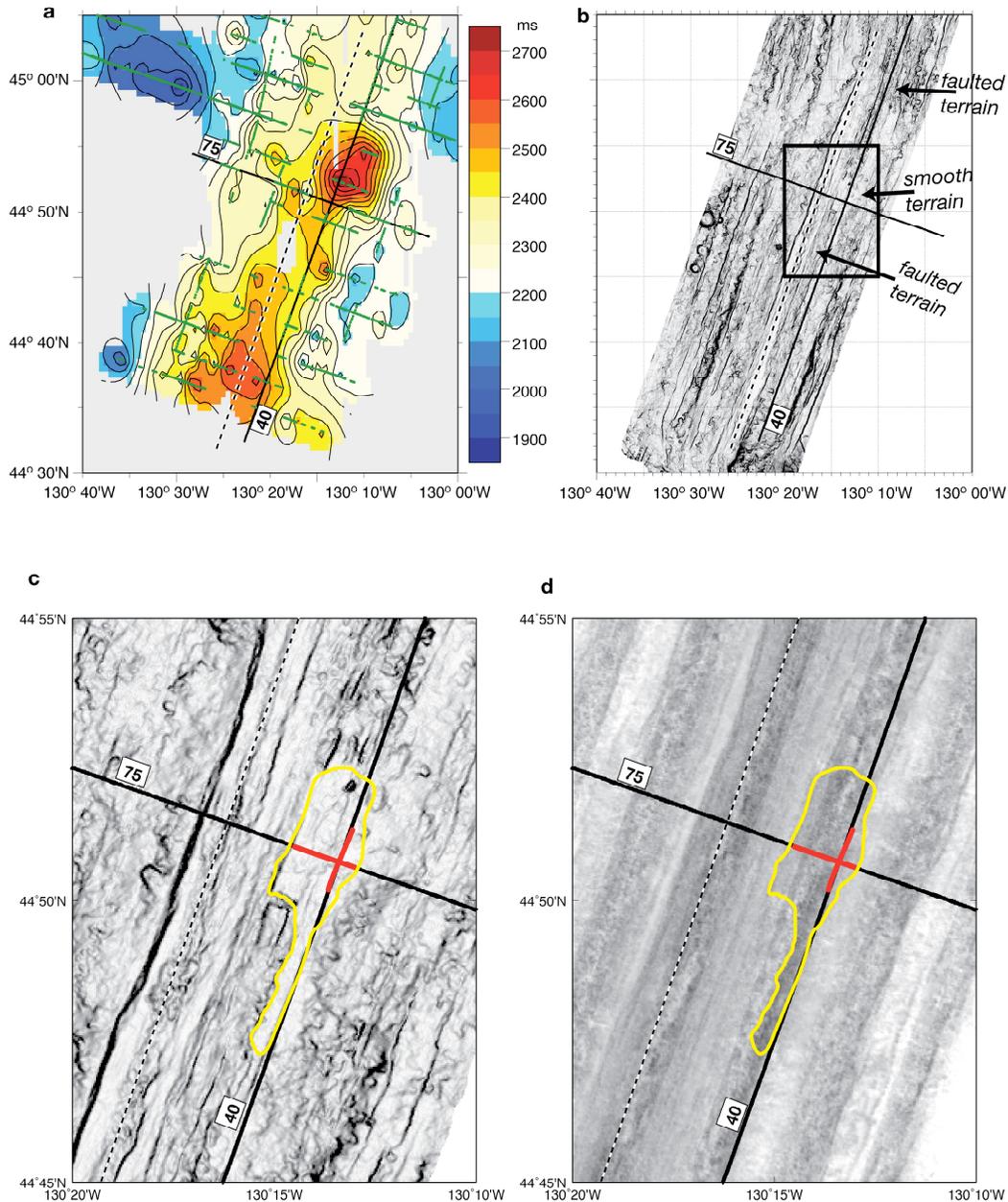
In the second model, the lower medium represents ultramafic rocks. We use the following elastic parameters:  $v_p = 7,700 \text{ m s}^{-1}$ ,  $v_s = 4,400 \text{ m s}^{-1}$  and

$\rho = 3,200 \text{ kg m}^{-3}$ . These parameters represent average values for wehrlite and harzburgite compositions<sup>37</sup> after correcting for the expected elevated temperatures of the near-axis lower crust (900 °C; ref. 7), assuming that  $v_p$  decreases with temperature as  $-0.54 \times 10^{-3} \text{ km s}^{-1} \text{ K}^{-1}$  (ref. 35) and a Poisson's ratio of 0.26 (ref. 36).

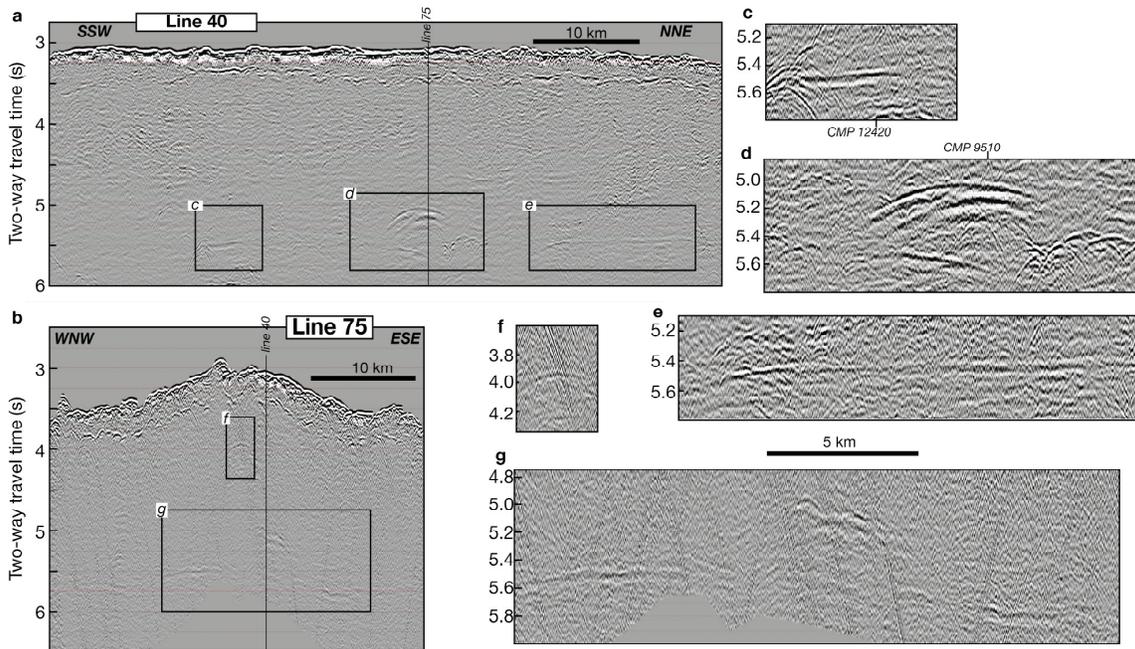
The predicted reflection coefficients do not exactly reproduce the observed AVO responses because the simplified models do not include effects such as energy attenuation and cannot fully describe complex reflectivity. True reflection coefficients for the LCML and Moho could not be estimated from the sea-floor multiple-to-primary amplitude ratio due to sea-floor diffractions that interfere with the LCML at zero offset.

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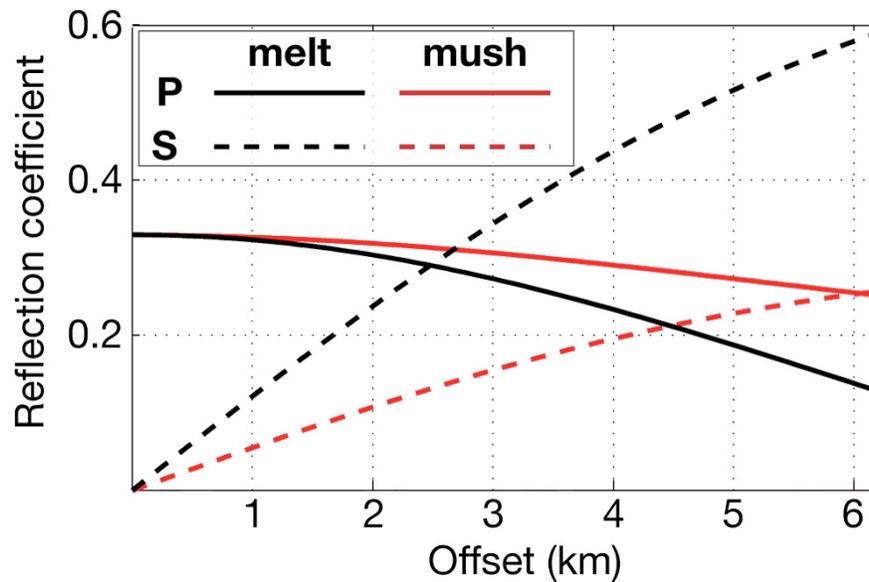
## SUPPLEMENTARY INFORMATION



**Figure S1.** **a.** Map of seafloor-to-Moho two-way traveltime (color scale in milliseconds to the right) derived from all of the EW0207 seismic profiles available in the area<sup>21</sup>. Green dots are locations where Moho was identified and picked; data points were interpolated onto a regular grid using a minimum-curvature algorithm. Contours are every 50 ms. Black lines are seismic profiles 75 and 40; dashed line is the ridge axis. **b.** Seafloor slope map of Cleft segment derived from high-resolution EM-300 multibeam bathymetry<sup>30</sup>. Light gray shading corresponds to sub-horizontal terrain and dark shading to slopes higher than 20 degrees. Box corresponds to area shown in **b** and **c**. Other lines are as in **(a)**. Details of seafloor slope **(c)** and EM-300 seafloor backscatter (i.e., reflectivity) **(d)** near the intersection of profiles 75 and 40 (black lines). Red lines indicate the plan-view extent of the LCML reflection event shown in Figs. 2 and 4. Yellow line encloses area of the seafloor where seafloor reflectivity is higher (darker) than the surrounding seafloor.



**Figure S2.** Unmigrated stacks of reflection profiles shown in Fig. 2 without interpretation. Processing steps are indicated in the Methods section.



**Figure S3.** Predicted reflection coefficients versus source-receiver offset for  $P_{LCML}P$  (solid lines) and  $P_{LCML}S$  (dashed lines) for the “melt” model described in Fig. 3 and Methods section (black lines), and a partially-crystallized sill model (“mush”, red lines:  $V_P=3400 \text{ m s}^{-1}$ ,  $V_S=3200 \text{ m s}^{-1}$ ,  $\rho=2700 \text{ kg m}^{-3}$ ). The  $P_{LCML}S$  is predicted to have maximum amplitudes at offsets of 10–12 km (not shown in here), beyond the maximum offset of our data (6.2 km). Horizontal axis is source-receiver offset.