## Iceberg discharges of the last glacial period driven by oceanic circulation changes

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Proxy data reveal the existence of episodes of increased deposition of ice-rafted detritus in the North Atlantic Ocean during the last glacial period interpreted as massive iceberg discharges from the Laurentide Ice Sheet. Although these have long been attributed to self-sustained ice sheet oscillations, growing evidence of the crucial role that the ocean plays both for past and future behavior of the cryosphere suggests a climatic control of these ice surges. Here, we present simulations of the last glacial period carried out with a hybrid ice sheet-ice shelf model forced by an oceanic warming index derived from proxy data that accounts for the impact of past ocean circulation changes on ocean temperatures. The model generates a time series of iceberg discharge that closely agrees with ice-rafted debris records over the past 80 ka, indicating that oceanic circulation variations were responsible for the enigmatic ice purges of the last ice age.

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ompared with the present interglacial period, the last glacial period (LGP) (~110–10 ka before the present), and almost certainly previous ones (1), were characterized by substantial climatic variability on millennial timescales. This variability is mainly manifested in two types of events. Dansgaard-Oeschger (D/O) events are most notable in Greenland ice core records and involve decadal-scale warming of more than 10 K (interstadials) followed by slow cooling lasting several centuries and a final more rapid fall to cold background (stadial) conditions (2). Heinrich (H) events consist of massive iceberg discharges from the Laurentide Ice Sheet at intervals of  $\sim$ 7 ka during peak glacial conditions throughout the LGP (3). Both D/O and H events are associated with widespread centennial- to millennial-scale climatic changes, including a synchronous temperature response over the North Atlantic and an antiphase temperature relationship over Antarctica and most of the Southern Ocean, as revealed by a wealth of deep-sea sediments, ice core, and terrestrial records (4). The Atlantic meridional overturning circulation (AMOC) is thought to play a central role in these abrupt glacial climatic changes. Although the paleoceanographic evidence on this link is scarce and mostly restricted to a few highresolution deep-sea sediment records of the last deglaciation (5, 6), both modeling studies and reconstructions provide strong support for the hypothesis that D/O events were caused by reorganizations of the AMOC (7, 8). H events, identified as enhanced ice-rafted detritus (IRD) in North Atlantic deep-sea sediments (3, 9), occur during climatic minima of the Northern Hemisphere. They have classically been attributed to internal oscillations of the Laurentide (10) and assumed to lead to important disruptions of the Atlantic Ocean circulation (11). However, paleoclimate data have revealed that most H events likely occurred about a thousand years after North Atlantic Deep Water (NADW) formation had already slowed down or largely collapsed (12, 13), implying that the initial AMOC reduction could not have been caused by the H events themselves. This evidence directly conflicts with the common interpretation that

freshwater fluxes representing the iceberg discharges caused the shift into cold (i.e., stadial) conditions. This furthermore highlights the need for a new paradigm through which to understand the triggering mechanism of H events. As already advanced one decade ago (14), any new theory should be able to account for the fact that the cold periods in which H events appear are not caused by the iceberg discharges and that the latter occur systematically several centuries after the North Atlantic cooling. More recently, the interaction between ocean circulation and ice sheet dynamics has been suggested to play a major role in triggering H events (15-17). This hypothesis has been assessed in particular for the first H event, H1, with both models and data showing that reduced NADW formation and a weakened AMOC lead to subsurface warming in the Nordic and Labrador Seas. This results in rapid melting of the Labrador ice shelves causing substantial ice stream acceleration and enhanced iceberg discharge (18–20).

Here, we investigate the effects of oceanic circulation changes associated with millennial-scale climate variability on the Laurentide Ice Sheet dynamics within a more realistic modeling framework. To this end, we drive a hybrid ice sheet–ice shelf model (21) with time-varying oceanic subsurface temperature fields for the LGP (*Materials and Methods*) obtained by combining glacial climate simulations and information from proxy data. Climatic boundary conditions are otherwise fixed to glacial conditions, so that the only external forcing felt by the ice sheet model is the change in subsurface ocean temperatures. These are translated into basal melting rates via a linear equation

## **Significance**

Periodic episodes of massive iceberg discharges from the large Northern Hemispheric ice sheets into the North Atlantic Ocean occurred throughout the last glacial cycle. It is still not clear whether they resulted from internal ice dynamics alone or were possibly externally driven. Results of our simulations of the Laurentide Ice Sheet forced by oceanic circulation changes support the hypothesis that these ice discharges were induced by the collapse of a buttressing ice shelf and the subsequent acceleration of inland ice streams. This provides a new basis for understanding the dynamics of the coupled cryosphere–climate system of glacial cycles. Additionally, it has strong implications for the stability of the marine parts of the Antarctic ice sheet given anthropogenic oceanic warming.

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dependent on a single tunable parameter (see SI Text for details and sensitivity tests). Climate simulations are performed with a global atmosphere-ocean model for glacial stadial and interstadial states (with weak and strong AMOC states, respectively) (22). These provide the range and spatial distribution of oceanic temperatures felt by the ice sheet. The temporal millennial-scale variability is based on a proxy-derived index used to interpolate in time between the stadial and interstadial ocean temperature fields. To produce this index, we assume that millennial-scale variability registered in the Greenland Ice Core Project (GRIP)  $\delta^{18}$ O ice core record (2) reflects variations in the North Atlantic oceanic state (Fig. 1A). To characterize the latter, following previous work (1), we use a threshold in the derivative of the  $\delta^{18}$ O GRIP signal to determine the timing of stadial to interstadial transitions (Fig. 1B). This allows for an objective classification of climatic states into stadials or interstadials (i.e., cold and warm surface periods). We furthermore assume millennial-scale variability as registered in the GRIP record reflects variations in NADW formation that have an imprint on subsurface temperatures in antiphase with respect to the surface state. Stadials are thus associated with periods of reduced NADW formation and weakened AMOC and warm subsurface temperatures, whereas during interstadials a stronger AMOC with active NADW formation cools the subsurface, in agreement with previous studies (15, 16, 19, 20, 23, 24). Considering a fast relaxation time of the subsurface temperature when convection resumes, and slow relaxation when convection is weak (20), we generate a subsurface warming index that slowly peaks during stadial climatic periods and more abruptly collapses when entering interstadial climatic periods (Fig. 1C). This index is thus directly derived from the GRIP time series and represents the only external forcing to the ice sheet model (see Materials and Methods and SI Text for details).

The application of subsurface oceanic forcing to the ice sheet model induces significant millennial-scale variability in the otherwise stable Laurentide Ice Sheet, as reflected in the velocity at the Hudson Strait outlet and iceberg discharge into the ocean (Fig. 2). For almost every peak in subsurface warming, there is a corresponding large and abrupt acceleration of the ice flow. Transitions between slow and fast states of the Hudson Strait ice stream occur several times during the LGP, with velocities varying between  $\sim 1,000 \text{ m} \cdot a^{-1}$  during buttressing periods and ~4,000 m·a<sup>-1</sup> during periods of ice shelf breakup. The magnitude of the velocity does not directly correlate with the magnitude and duration of the subsurface warming, however, because of the competing timescales of ice sheet growth, ice advection from inland, and ice shelf breakup and growth. These three mechanisms lead to a nonlinear response of the system that appears to modulate the dynamics of the floating and inland ice in this region. When interstadial subsurface (i.e., cold) oceanic conditions are applied, the Labrador Sea ice shelf experiences low melt rates and can extend far enough to reach the western coast of Greenland (Fig. 3). In this way, significant backforce is felt by the Hudson Strait ice stream and velocities are greatly reduced. This allows the main Laurentide ice dome to grow and subsequently advect ice from inland toward the margin because of a permanently active Hudson Strait ice stream, preconditioning the ice sheet for more ice discharge into the ocean. When stadial subsurface (i.e., warm) oceanic conditions are applied, the ice shelf melts away from Greenland and no longer buttresses the ice stream that feeds it (Fig. 3). This allows a surge of velocity at the mouth of the ice stream, which propagates inland over several centuries and results in a significant increase in ice discharge into the Labrador Sea (Fig. 4). The magnitude of such a discharge event depends on the state of the ice sheet before the ice shelf collapse.



**Fig. 1.** Derivation of the Labrador Sea subsurface oceanic index. (A) GRIP d18O ice core record. (B) Smooth derivative of GRIP d18O record (red) with positive and negative thresholds which define transitions between stadial and interstadial states (grey). (C) Subsurface warming index. The cold subsurface state corresponds to an interstadial state (i.e., warm surface "climatic" state) with a mean subsurface (700- to 1,100-m depth) temperature of -0.9 °C. The warm subsurface state corresponds to a stadial state (i.e., cold surface "climatic" state) with a mean subsurface temperature of 1.1 °C.



Fig. 2. (A) Labrador Sea subsurface oceanic index; (B) simulated Hudson Strait ice velocity (in kilometers per year); (C) simulated Labrador Sea calving rate (in Sverdrups); (D) magnetic susceptibility from core MD95-2024 (45.7°W, 50.2°N) (25); (E) lithic fraction from core JPC-13 (33.5°W, 53.1°N) (26). For the comparison, the timescales of the proxy data were converted to the SS09 timescale of the GRIP record (2).

The simulated time series of calving into the Labrador Sea compares very well with proxies of calving obtained from marine sediment cores from the North Atlantic (Figs. 2 and 3). Both a high-resolution record of magnetic susceptibility from core MD95-2024 (45.7°W, 50.2°N) (25) and a record of lithic fraction from core JPC-13 (33.5°W, 53.1°N) (26), i.e., IRD proxies, show the same timing of peaks corresponding to major discharge events. In some isolated cases, such as between H4 and H3, or between H2 and H1, the simulated time series agrees better with the latter core. However, no spurious discharge events are simulated that are not apparent in at least one core. A comparison of the time series of the prescribed subsurface warming, calving, and IRD proxies highlights the fact that, for every strong peak in calving (i.e., H event), there is a necessary peak in subsurface warming. Several sensitivity tests with Grenoble ice shelf and land ice model (GRISLI) show that the amount of calving produced is a result of the nonlinear preconditioning of the ice sheet (SI Text). However, in our simulations the triggering mechanism for large ice discharges is always an ice shelf breakup precipitated by subsurface warming in the Labrador Sea. The strongest calving events are furthermore found to take place during the longest subsurface warming periods. We would expect the same relationship if subsurface temperature reconstructions were available, reflecting positive feedbacks operating between NADW formation and ice discharge (16). A more persistent reduction in NADW formation results in longer periods of subsurface warming. This in turn has a larger impact on ice shelves, which tends to increase ice discharge and suppress NADW formation further.

H events are among the most dramatic examples of millennialscale variability of the Quaternary climate and their interpretation has remained elusive for decades. In recent years, the increasing availability of observations of the present-day ice sheets has confirmed the unexpected and crucial role that the ocean exerts on the dynamics of the ice sheets (27). Ice shelves represent the necessary interface for this coupled system. Whereas little information exists for a constrained reconstruction of the floating parts of the Laurentide, the maximum extent of the ice shelf simulated here is restricted to the continental shelf area between Greenland and the Hudson Strait. Such a configuration does not appear to contradict the relatively sparse proxy data available in this region (28), and is glaciologically consistent. Furthermore, the existence of a persistent ice stream through the Hudson Strait, as simulated here, is supported by geological evidence and modeling (29, 30).

Combined with the simulations presented here, the fact that the subsurface warming index generated from GRIP  $\delta^{18}$ O data



Fig. 3. Laurentide ice stream velocities (in kilometers per year) before (*Left*) and during (*Right*) H event 2, along with locations of the cores of the IRD proxies shown in Fig. 2. The dashed line in the right panel indicates the location of the profiles shown in Fig. 4.

aligns so well with the IRD proxies lends strong support to the hypothesis that millennial-scale glacial ice discharges are the result of a response to oceanic forcing. A characteristic time longer than the forcing timescale is the result of the non-linearities of the ice sheet/ice shelf system. These arise from the different characteristic times of the ice shelf breakup and regrowth and by the time needed by the ice sheet to propagate the signal from its oceanic perturbation across the ice streams (Fig. 4). These phenomena favor the occurrence of resonance in the system and finally determine the observed pacing of  $\sim$ 7 ka.

Our simulations provide a physically based framework through which to understand the coupled ice sheet–ocean system. Open questions remain concerning the relationship between IRD proxies and actual calving rates, which can result from outburst floods, iceberg melting, and ocean circulation changes (31). One important related aspect concerns the fact that it is very difficult to constrain the melting rates that icebergs experience during their trip across the North Atlantic. This allows for alternative explanations considering the observed IRD belt as mainly the reflection of colder oceanic temperatures when Heinrich layers were formed (32). Under this interpretation, however, the amount of IRDs in marine cores close to the ice sheet source would reflect a signal absent of Heinrich-like events. This seems not to be the case, because Heinrich peaks can be observed in cores of the Labrador Sea (33). However, the explanation for the ultimate causes behind the underlying glacial oceanic variability remains elusive. Nonetheless, the work presented here shows that proxies and modeling reveal a consistent picture of the origin of the massive iceberg discharges of the last glacial cycle, including the enigmatic H events.

## **Materials and Methods**

The ice sheet model GRISLI simulates the 3D evolution of the Laurentide using a hybrid ice sheet/ice shelf approach. GRISLI is one of the few models able to properly deal with both grounded and floating ice on the paleo-hemispheric scale, because it explicitly calculates grounding line migration, ice stream velocities, and ice shelf behavior. Inland ice deforms according to the stress balance using the shallow ice approximation (34, 35). Ice shelves are described following ref. 36, and ice streams (areas of fast flow, typically larger than  $\sim 10^2 \text{ m} \cdot \text{y}^{-1}$ ) are considered as dragging ice shelves, allowing basal movement of the ice (37). Basal drag under ice streams is proportional to ice velocity and to the effective pressure. The locations of the ice streams are determined by the basal water within areas where the sediment layer is saturated. Contrary to the classic "binge-purge" theory (10), basal ice movement is computed here under the shallow shelf approximation. Rapid ice flow areas are therefore simulated in a more realistic dynamical approach (37). As a consequence, internal basal temperature oscillations, and thereby Laurentide instabilities, are found to vanish. In the absence of any oceanic forcing, the Laurentide Ice Sheet reaches a nonoscillatory steady



Fig. 4. Ice sheet profiles of the Laurentide (as indicated in Fig. 3) before and during HE2. Time series show ice-shelf thickness (gray; in meters), basal stress (10<sup>5</sup> Pa), velocity (in kilometers per year), and thickness (in meters) for the upstream (magenta) and downstream (dark blue) sections of the Hudson Strait ice stream. The background shading in the right panel represents buttressed (light blue), transition (light red) and unbuttressed (white) periods.

state. Climate simulations are performed with CLIMBER-3 $\alpha$ , which includes an oceanic general circulation model.

Basal melting rates under the ice shelves are computed here using a linear relationship on the difference between the subsurface temperature,  $T_o$ , and the temperature of the freezing point of salty waters,  $T_f$ :

$$B = \kappa (T_o - T_f),$$
<sup>[1]</sup>

where *B* is the basal melt rate under the Labrador Sea floating ice (in meters per year).  $T_o$  is the subsurface temperature of the Labrador Sea and its evolution through time is given by the following:

$$T_o \equiv T_o(t) = T_{is} + \alpha(t) \cdot (T_s - T_{is}).$$
 [2]

Then,

$$\min(T_o) = T_{is}, \quad \text{when} \quad \alpha = 0$$
$$\max(T_o) = T_s, \quad \text{when} \quad \alpha = 1,$$

where  $\alpha(t)$  is the subsurface warming index shown in Fig. 1, and  $T_s$  and  $T_{is}$  are the mean Labrador Sea subsurface temperatures for a stadial and an interstadial period, respectively.

The time series of IRD content from proxy data were converted to a common timescale with the model forcing for more direct comparison of

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the results. Namely, we used the original SS09 timescale (38) of the GRIP dataset (39) for all time series. The IRD record of ref. 25 was originally provided on the GISP2 timescale (40). Conveniently, the GRIP dataset included equivalent times between the SS09 and GISP2 timescales, allowing direct conversion of this time series to the SS09 timescale via linear interpolation. The IRD record of ref. 26 was provided on the SFCP timescale (41). Here, the equivalent times were only available for the newer SS09sea (42) timescale. In this case, we compared GRIP  $\delta^{18}$ O values available on the SFCP timescale with the original data on the SS09 timescale, and optimized for the time corrections at 14 tie points (with linear interpolation in between) that would make the former match the latter. This procedure is accurate enough to allow for visual comparison on millennial timescales. For example, it is able to reproduce the SS09  $^{18}$ O values with a root mean square error of 0.1 per mille. In Dataset S1, we provide equivalent times for the last 80 ka.

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