# **Anisotropic Internal Stress in Landfast Ice from Canadian Arctic Archipelago**

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## **1. Introduction and Description**

A better understanding of landfast ice is indispensable to improve the simulation of high latitude climate since landfast ice covers  $1.8 \times 10^6$  km<sup>2</sup> (~ 20 %) of the Arctic Ocean and its average annual duration is over seven months. An internal Ice Stress Buoy (ISB) was deployed near the center of a multiyear floe in the Viscount Melville Sound in the Canadian Arctic Archipelago (CAA). The buoy collected hourly internal sea-ice stress, GPS coordinates, air temperature and orientation between October 10, 2010 and August 17, 2011, which include irresponsible stress data from June 19 - August 17, 2011.

## 2. Ice Motion and Landfast Season

## **3. Firstyear Sea-Ice Strength Estimates**





Figure: The time series of speed of the multiyear ice drift from October 10, 2010 to August 17, 2011.

The ISB was deployed on October 10, 2010 (Julian Day 283) at 7415' North, 10805' West in the Viscount Melville Sound within the CAA. From January 18 (JD 383) to July 10 (JD 556) the ice was landfast.



Figure: Time series of floe speed (in blue) and GEM reanalysis wind data [5] (in red) before and after sea ice became landfast on January 18 (JD 383)

Integrating the wind force acting on ice floe by [2]:

$$au_a imes L$$
 (1)

Figure: The drift of the multiyear ice floe from October 10, 2010 to August 17, 2011.

### 4. Internal Ice Stresses

Sudden Short-Lived Peaks in Stresses
The internal sea-ice stress data shows
 high frequency fluctuations on
 November 19 - 23 (JD 323 - 327),
 December 4 - 5 (JD 338 - 339), December
 14 - 15 (JD 348 - 349) and on January 2 - 6
 (JD 366 - 371) when the floe speed drops
 to zero momentarily We interpret these
 sudden short-lived peaks in stresses
 accompanied by floe stoppage as
 collisions between multiyear ice floes.



**Figure**: Time series of floe speed (in blue) and stress data ( $\sigma_2 - \sigma_1$ : in green).

where  $\tau_a$  ( =  $C_{da}\rho_a |\mathbf{u}_a|^2$  ) is surface wind force, and *L* (170 km) is the full width of the Parry Channel

Estimate firstyear ice strength is around 30 kN/m.

#### ▶ $\sigma_2 - \sigma_1$ and wind speed

 $\sigma_2 - \sigma_1$  is often interpreted as dynamic stresses [1], however,  $\sigma_2 - \sigma_1$  has very low correlation with wind speed ( $r^2 = 0.13$ ) in CAA. We interpret this result as an indication of the presence of anisotropy in coastal sea ice of the CAA by the preferred c-axis alignment [4].



Figure: The time series of (top)  $\sigma_2 - \sigma_1$  and (bottom) wind speed around the floe [5].

## Principal Stresses and Air Temperature

Both  $\sigma_1$  and  $\sigma_2$  has high correlation with temperature (r<sup>2</sup> = 0.81 and 0.67, respectively).



Figure: The time series of (top) principal stresses ( $\sigma_1$  and  $\sigma_2$ ) and (bottom) air and sensor temperatures between October 10, 2010 to June 17, 2011. Note that principal stresses are plotted negative upward to see clearer correlation with surface air temperature.

## 5. Anisotropic Thermal Stress Model

▶ Model Equations 1: A 1.5D anisotropic thermal stress model based on linear elasticity theory

$$\sigma_{11}^{th}(z) = \frac{E_1}{1 - \nu_1 \nu_2} \left( \left[ \frac{\int_{-h}^0 E_1 \alpha \Delta \theta_{(z)} dz}{\int_{-h}^0 E_1 dz} - \alpha \Delta \theta_{(z)} \right] + \nu_2 \left[ \frac{\int_{-h}^0 E_2 \alpha \Delta \theta_{(z)} dz}{\int_{-h}^0 E_2 dz} - \alpha \Delta \theta_{(z)} \right] \right)$$
  
$$\sigma_{22}^{th}(z) = \frac{E_2}{1 - \nu_1 \nu_2} \left( \left[ \frac{\int_{-h}^0 E_2 \alpha \Delta \theta_{(z)} dz}{\int_{-h}^0 E_2 dz} - \alpha \Delta \theta_{(z)} \right] + \nu_1 \left[ \frac{\int_{-h}^0 E_1 \alpha \Delta \theta_{(z)} dz}{\int_{-h}^0 E_1 dz} - \alpha \Delta \theta_{(z)} \right] \right)$$

where internal ice temperature is calculated by the 1D thermodynamic model of [3].

- Youngs modulus,  $E_1 = 0.2$  GPa (parallel to c-axis) and  $E_2 = 0.14$  GPa (normal to c-axis) are required.
- The simulation is in some agreement with observed stress data.
- The principal stresses switch order when large surface warming and cooling are present.
- The deviation between the measured and simulated stress grows as the surface temperature gets colder. → due to stress relaxation via viscous creep.
- Model Equations 2: Adding a simple relaxation term



• The relaxation time constant = 110 days is



Figure: Time series of observed principal stresses and simulated internal sea ice thermal stresses at the depth of the sensor. No snow cover and no stress relaxation are assumed in the simulation.



Best Match



- required.
- The simulation is in better agreement with observed stress data.
- The stress relaxation via viscous creep present. → The compressive stress would present at the beginning of the melt season instead of zero stress.

Figure: Time series of observed principal stresses and simulated internal sea ice thermal stresses at the depth of the sensor. No snow cover and stress relaxation via viscous creep with a relaxation time scale of 110 days (2640 hours) are assumed in the simulation.

Figure: The time series of (top) stress vertical distribution and (bottom) stress at the sensor depth.

- The simulation is in good agreement with observed stress data.
- The relaxation time constant decreases to 720 hours.
- Compressive stress at the beginning of melt season is assumed. -30 kPa in  $\sigma_1$  and -50 kPa in  $\sigma_2$  are required for best match.
- The vertical distribution of the thermal stress from the simulation shows that when tensile stresses develop near ice surface, compressive stresses develop near the bottom of the ice.

## 7. References

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# 6. Conclusion

- ► Landfast season in the Viscount Melville Sound last nearly six months from January 18 to July 10, 2011.
- ► A few events when dynamic stresses are large associating with floe-floe collision.
- ► A compressive strength of firstyear sea ice of 30 kN/m is estimated
- The observed stress data is highly correlate with changes in surface air temperature and not significantly correlated with the magnitude of the wind speed.
- Anisotropy in sea ice is responsible for the difference in principal stresses as opposed to dynamic stresses in region where sea ice is isotropic [e.g.][]Richter-Menge2002, Hutchings2010.
- A 1.5D anisotropic thermal stress model for sea ice including the effect of stress relaxation via viscous creep results in a good agreement with the observed principal stresses, when stress relaxation via viscous creep is considered.